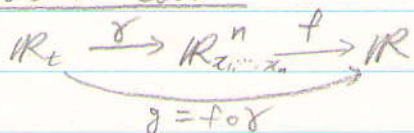


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Chain Law



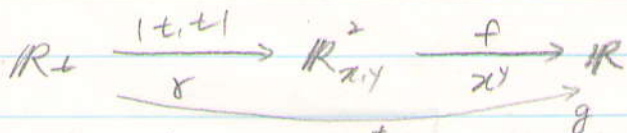
$$\gamma = \begin{pmatrix} \gamma_1 \\ \vdots \\ \gamma_n \end{pmatrix}$$

$$g(t+\epsilon) = f(\gamma(t+\epsilon)) \\ = f(\gamma_1(t+\epsilon), \dots, \gamma_n(t+\epsilon))$$

$$= f_1 \gamma_1 \epsilon + f_2 \gamma_2 \epsilon + \dots$$

change due to change in 1st coord.

$$= \left(\sum \frac{\partial f}{\partial x_i} \gamma_i'(t) \right) \epsilon$$

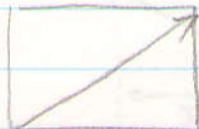


$$\gamma(t) = \begin{pmatrix} t \\ t \end{pmatrix}, f(x,y) = x^y$$

$$g = f \circ \gamma = t^t$$

$$\gamma_1' = 1, \gamma_2' = 1, \gamma_1' = 1, \gamma_2' = 1$$

$$g'(t) = \frac{\partial f}{\partial x} \cdot 1 + \frac{\partial f}{\partial y} \cdot 1$$



$$= \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y}$$

Problem

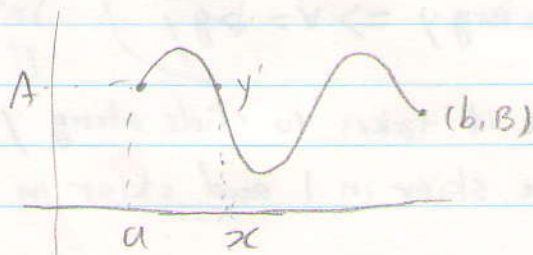
Minimize

$$F: \mathbb{R}^3 \rightarrow \mathbb{R}$$

sufficiently differentiable

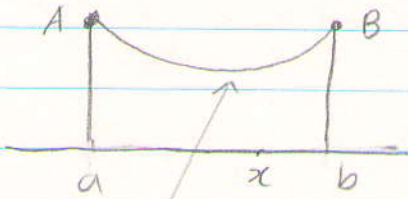
$$J(y) = \int_a^b F(x, y, y') dx$$

Among all sufficiently differentiable functions $y: [a, b] \rightarrow \mathbb{R}$
 s.t. $y(a) = A$ and $y(b) = B$



examples

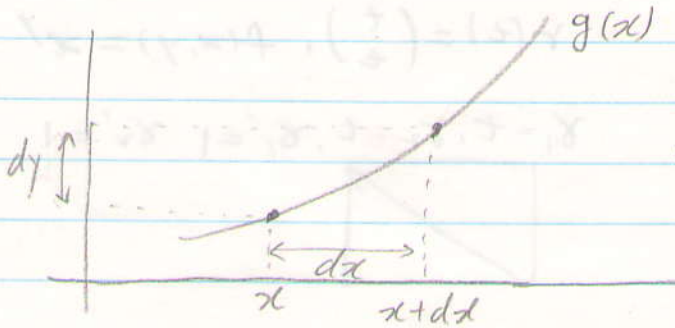
① The "power line" problem



$$y = \cosh x = \frac{e^x + e^{-x}}{2}$$

$J(y)$ = potential energy of y
 the potential energy of a tiny piece of string is proportional to its height, and its mass

$$J(y) = \int_a^b y \left(\begin{array}{l} \text{mass of} \\ \text{string} \\ \text{between} \\ x \text{ and } x+dx \end{array} \right) \propto \int_a^b y \left(\begin{array}{l} \text{length of} \\ y \text{ between} \\ x \text{ and } x+dx \end{array} \right) = \int_a^b y \cdot \sqrt{1+(y')^2} dx$$



$$\begin{aligned} \text{length} &= \sqrt{(dx)^2 + (dy)^2} \\ &= \sqrt{1 + \left(\frac{dy}{dx}\right)^2} dx \\ &\rightsquigarrow \sqrt{1+(y')^2} dx \end{aligned}$$

$$F(x_1, x_2, x_3) = x_2 \sqrt{1+x_3^2}$$

② $L(q) = \int \frac{1}{2} m \dot{q}^2 - V(q(t))$

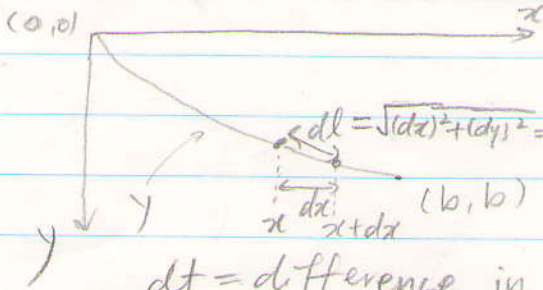
"J" "y" "x"

kinetic energy potential energy at q

$V: \mathbb{R} \rightarrow \mathbb{R}$

q : position of particle as function of time.

③ The Brachistochrone



$$\frac{1}{2} m v^2 = m g y \Rightarrow v = \sqrt{2 g y}$$

$$\int_0^b \frac{\sqrt{1+(y')^2}}{\sqrt{2 g y}} dx$$

||

$$\int_0^b dt$$

||

$T(y)$ = time it takes to slide along y

dt = difference in time between skier in 1 and skier in 2

$$= \frac{dl}{v} = \frac{\sqrt{1+(y')^2}}{\sqrt{2 g y}} dx$$

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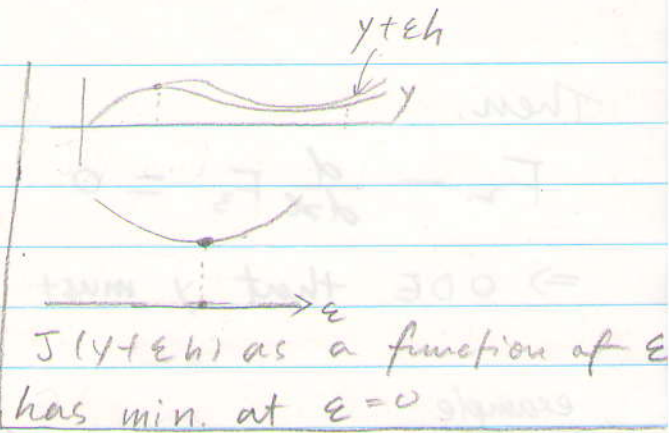
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Euler - Lagrange Equation

y is an extremum

$\Rightarrow \forall h, h: [a, b] \rightarrow \mathbb{R}, h(a)=0, h(b)=0$

$\frac{d}{d\epsilon} J(y + \epsilon h) \Big|_{\epsilon=0} = 0$



$\frac{d}{d\epsilon} J(y + \epsilon h) \Big|_{\epsilon=0} = \frac{d}{d\epsilon} \int_a^b F(x, y + \epsilon h, y' + \epsilon h') dx \Big|_{\epsilon=0}$

$= \int_a^b \frac{d}{d\epsilon} F(x, y + \epsilon h, y' + \epsilon h') dx \Big|_{\epsilon=0}$

$= \int_a^b (F_1 \cdot 0 + F_2 \cdot h + F_3 \cdot h') dx \Big|_{\epsilon=0}$

$F_2(x, y + \epsilon h, y' + \epsilon h')$

$= \int_a^b (F_2(x, y + y')h + F_3(\dots)h') dx$

integrate by parts

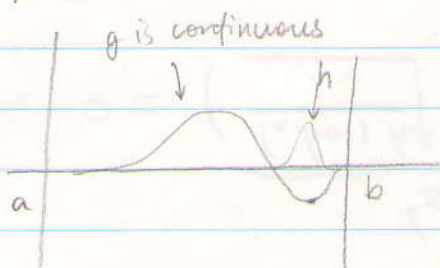
$= \int_a^b (F_2 h - (\frac{d}{dx} F_3) h) dx + \text{boundary terms } F_3 \cdot h \Big|_a^b \frac{h(a)=0}{h(b)=0} = 0$

$= \int_a^b \underbrace{(F_2 - \frac{d}{dx} F_3)}_{g''(x)} \cdot h(x) dx \text{ (should be } 0 \forall h)$

$\Rightarrow \forall h, \int_a^b g \cdot h = 0$

$\Rightarrow g \equiv 0$

proof



Then,

$$F_2 - \frac{d}{dx} F_3 = 0 \quad (\text{Euler-Lagrange Equation})$$

\Rightarrow ODE that y must satisfy

example

$$J = \int_a^b \underbrace{\left[\frac{1}{2} m (\dot{y})^2 - V(y) \right]}_{F(x, y, \dot{y})} dx$$

$$F(u_1, u_2, u_3) = \frac{1}{2} m u_3^2 - V u_2$$

$$F_2 = -V'(u_2) \quad \text{and} \quad F_3 = m u_3$$

$$\text{E-L: } F_2 - \frac{d}{dx} F_3 = 0$$

$$\Rightarrow -V'(y) - \frac{d}{dx} (m\dot{y}) = 0$$

$$m\ddot{y} = -V'(y)$$

$$ma = F$$

same equation

note

$$F_y - \frac{d}{dx} F_{y'} = 0 \quad (\text{Euler-Lagrange Equation})$$

example (Brachistochrone)

$$F = \sqrt{\frac{1+y'^2}{y}}$$

$$\text{E-L: } \underbrace{-\frac{1}{2} \sqrt{\frac{1+y'^2}{y^3}}}_{F_y} - \frac{d}{dx} \left(\underbrace{2y' \cdot \frac{1}{2} \sqrt{\frac{1}{y(1+y'^2)}}}_{F_{y'}} \right) = 0$$