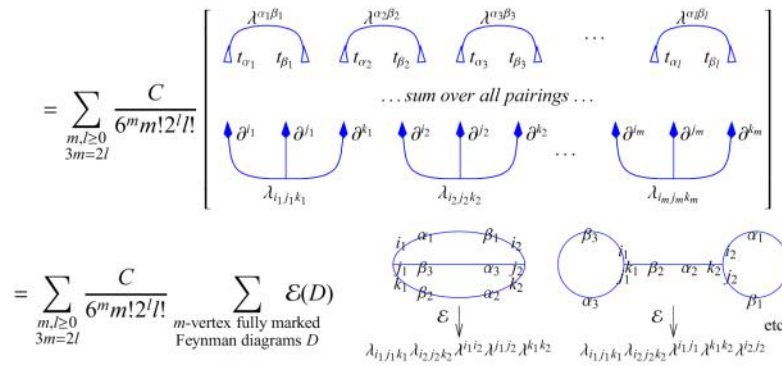


From Gaussian Integration to Feynman Diagrams

We wish to understand $\int_{A \in \Omega^1(\mathbb{R}^3, \mathfrak{g})} \mathcal{D}A \text{hol}_\gamma(A) \exp\left(\frac{ik}{4\pi} \int_{\mathbb{R}^3} \text{tr} A \wedge dA + \frac{2}{3} A \wedge A \wedge A\right)$.

As a warm up, suppose (λ_{ij}) is a symmetric positive definite matrix and (λ^{ij}) is its inverse. Denote by $(x^i)_{i=1}^n$ the coordinates of \mathbb{R}^n , let $(t_i)_{i=1}^n$ be a set of "dual" variables, and let ∂^i denote $\frac{\partial}{\partial t_i}$. Also let $C := \frac{(2\pi)^{n/2}}{\det(\lambda_{ij})}$. Then

$$\int_{\mathbb{R}^n} e^{-\frac{1}{2} \lambda_{ij} x^i x^j + \frac{1}{6} \lambda_{ijk} x^i x^j x^k} = C e^{\frac{1}{6} \lambda_{ijk} \partial^i \partial^j \partial^k} e^{\frac{1}{2} \lambda^{\alpha\beta} t_\alpha t_\beta} \Big|_{t_\alpha=0} = \sum_{\substack{m, l \geq 0 \\ 3m=2l}} \frac{C}{6^m m! 2^l l!} (\lambda_{ijk} \partial^i \partial^j \partial^k)^m (\lambda^{\alpha\beta} t_\alpha t_\beta)^l$$

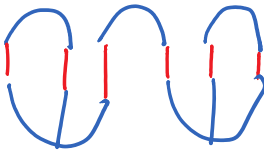


$$= C \sum_{\substack{m, l \geq 0 \\ 3m=2l}} \frac{C}{6^m m! 2^l l!} \sum_{\text{m-vertex fully marked Feynman diagrams } D} \mathcal{E}(D)$$

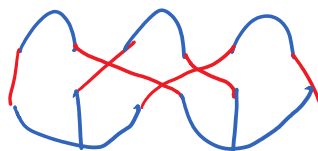
Claim. The number of pairings that produce a given unmarked Feynman diagram D is $\frac{6^m m! 2^l l!}{|\text{Aut}(D)|}$.

Proof of the Claim. The group $[(S_3)^m \times S_m] \times [(S_2)^l \times S_l]$ acts on the set of pairings, the action is transitive on the set of pairings P that produce a given D , and the stabilizer of any given P is $\text{Aut}(D)$. \square

Examples:



$|\text{Aut}(\text{O-O})| = 8$



$|\text{Aut}(\text{O})| = 12$

The Fourier Transform:

$(F: V \rightarrow \mathbb{C}) \Rightarrow (F: V^* \rightarrow \mathbb{C})$
 via $F(f) = \int_V f(v) e^{-i\langle v, \psi \rangle} dv$

Simple Facts:

1. $F(0) = \int_V f(v) dv$

2. $\frac{\partial}{\partial v_i} F \sim \sqrt{i} F$

3. $(e^{Q/2}) \sim e^{-Q'/2}$

where $Q'(v) = \langle v, L^{-1}v \rangle$

(That's the heart of the Fourier Inversion Formula.)

V : Vector space

dv : Lebesgue's measure on V .

Q : A quadratic form on V ;

$Q(v) = \langle L v, v \rangle$ where

$L: V \rightarrow V^*$ is linear

Constants $I = \int_V e^{\pm Q/2} dv$

$= \sum_{n=0}^{\infty} \frac{1}{n!} \int_V dv p^n e^{Q/2}$

$\sim \sum_{n=0}^{\infty} \frac{1}{n!} p^n(\partial_v) e^{-\frac{1}{2} Q'(v)} \Big|_{v=0}$

$= \sum_{n=0}^{\infty} \frac{\epsilon^{1/2}}{2^n n!} p^n(\partial) (Q')^n \Big|_{v=0}$

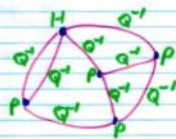
So $\int_V H(v) e^{\pm Q/2} dv$

$\sim H(\partial) e^{H(\partial)} e^{-Q'(v)/2} \Big|_{v=0}$

is



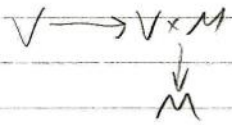
$= \sum_{\text{Diagrams}} c(D) (\text{Products of } Q's \text{ and one } H)$



Physics: "Physics should be gauge invariant"

Math: Could cure $\int_{Y_1} e^{i \int_{Y_2} A \wedge A}$ thanks to "invariance".

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$$A \in \Omega^1(M, \text{End } V)$$

si, pan $s: M \rightarrow V$ ρ

$$D_A s = ds + A \wedge s ; (D_A s)(\xi) = \xi s + A(\xi)(s)$$

$$g \circ D \mapsto D^g = g^{-1} D g ; g: M \rightarrow \text{Aut}(V) \text{ fibration}$$

$$A \mapsto g^{-1} d g + g^{-1} A g$$

$$v_0 \in V, \gamma: [0, 1] \rightarrow M$$

$$v \in p^{-1} \gamma: [0, 1] \rightarrow V$$

$$(\gamma^*(D_A))(\dot{\gamma}) = 0 \Rightarrow \left(\frac{d}{dt} v + A(\dot{\gamma}(t)) v \right) = 0$$

$$v(t); \dot{\gamma}(t) = e^{-\int_0^t A(\dot{\gamma}(s))} v_0$$

$$\gamma(t) = \left(I - \int_0^t A(\dot{\gamma}(s)) + \dots \right) v_0 \quad \frac{d}{dt} \text{hol}_\gamma(A)(t) = \text{hol}_\gamma(A)(t) A(\dot{\gamma}(t)) \text{hol}_\gamma(A)(t)^{-1}$$

$$\text{hol}_\gamma(A^g) = g \text{hol}_\gamma(A) g^{-1}$$

$$\int_M \text{tr}(\text{hol}_\gamma(A))$$

chern-simons

$$CS(A) = \int_M \text{tr}(A \wedge dA) + \frac{2}{3} A \wedge A \wedge A$$

$$CS(A) \in \mathbb{R}$$