

Cars, Interchanges, Traffic Counters, and some Pretty Darned Good Knot Invariants

More at oeβ/APAI



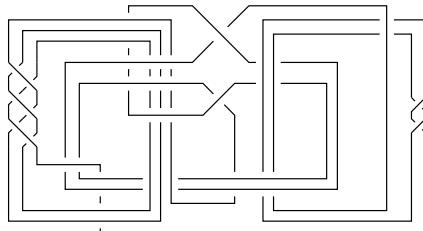
Abstract. Reporting on joint work with Roland van der Veen, I'll tell you some stories about ρ_1 , an easy to define, strong, fast to compute, homomorphic, and well-connected knot invariant. ρ_1 was first studied by Rozansky and Overbay [Ro1, Ro2, Ro3, Ov], it has far-reaching generalizations, it is dominated by the coloured Jones polynomial, and I wish I understood it. **Common misconception.** “Dominated” \Rightarrow “lesser”.



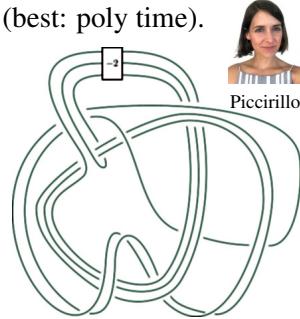
We seek strong, fast, homomorphic knot and tangle invariants.

Strong. Having a small “kernel”.

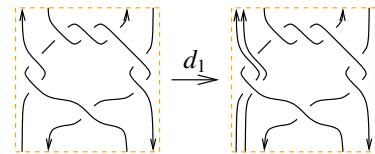
Fast. Computable even for large knots (best: poly time).



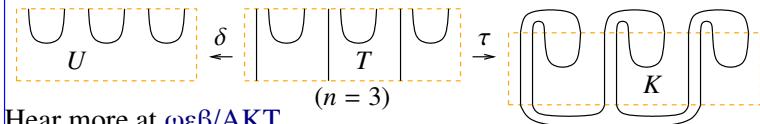
Gompf–Scharlemann–Thompson



Homomorphic. Extends to tangles and behaves under tangle operations; especially gluings and doublings:



Why care for “Homomorphic”? **Theorem.** A knot K is *ribbon* iff there exists a $2n$ -component tangle T with skeleton as below such that $\tau(T) = K$ and where $\delta(T) = U$ is the *untangle*:



Hear more at oeβ/AKT.

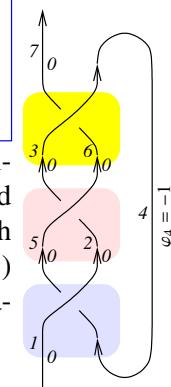
Acknowledgement. This work was partially supported by NSERC grant RGPIN-2018-04350 and by the Chu Family Foundation (NYC).

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Jones:

Formulas stay;
interpretations change with time.



Formulas. Draw an n -crossing knot K as on the right: all crossings face up, and the edges are marked with a running index $k \in \{1, \dots, 2n+1\}$ and with rotation numbers φ_k . Let A be the $(2n+1) \times (2n+1)$ matrix constructed by starting with the identity matrix I , and adding a 2×2 block for each crossing:

$$c : \begin{array}{cc} s = +1 & s = -1 \\ j+1 \uparrow & i+1 \uparrow \\ i & j \\ \downarrow & \downarrow \\ i & j \end{array} \rightarrow \begin{array}{c|cc} A & \text{col } i+1 & \text{col } j+1 \\ \hline \text{row } i & -T^s & T^s - 1 \\ \text{row } j & 0 & -1 \end{array}$$

Let $G = (g_{\alpha\beta}) = A^{-1}$. For the trefoil example, it is:

$$A = \left(\begin{array}{ccccccc} 1 & -T & 0 & 0 & T-1 & 0 & 0 \\ 0 & 1 & -1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -T & 0 & 0 & T-1 \\ 0 & 0 & 0 & 1 & -1 & 0 & 0 \\ 0 & 0 & T-1 & 0 & 1 & -T & 0 \\ 0 & 0 & 0 & 0 & 1 & -1 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{array} \right),$$

$$G = \left(\begin{array}{ccccccc} 1 & T & 1 & T & 1 & T & 1 \\ 0 & 1 & \frac{1}{T^2-T+1} & \frac{T}{T^2-T+1} & \frac{T}{T^2-T+1} & \frac{T^2}{T^2-T+1} & 1 \\ 0 & 0 & \frac{1}{T^2-T+1} & \frac{T}{T^2-T+1} & \frac{T}{T^2-T+1} & \frac{T^2}{T^2-T+1} & 1 \\ 0 & 0 & \frac{1-T}{T^2-T+1} & \frac{1}{T^2-T+1} & \frac{1}{T^2-T+1} & \frac{T}{T^2-T+1} & 1 \\ 0 & 0 & \frac{1-T}{T^2-T+1} & \frac{(T-1)T}{T^2-T+1} & \frac{1}{T^2-T+1} & \frac{T}{T^2-T+1} & 1 \\ 0 & 0 & 0 & \frac{T^2-T+1}{T^2-T+1} & 0 & 1 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{array} \right)$$

“The Green Function”



Wirtinger



Note. The Alexander polynomial Δ is given by

$$\Delta = T^{(-\varphi-w)/2} \det(A), \quad \text{with } \varphi = \sum_k \varphi_k, w = \sum_c s.$$

Classical Topologists: This is boring. Yawn.

Formulas, continued. Finally, set

$$R_1(c) := s(g_{ji}(g_{j+1,j} + g_{j,j+1} - g_{ij}) - g_{ii}(g_{j,j+1} - 1) - 1/2) \\ \rho_1 := \Delta^2 \left(\sum_c R_1(c) - \sum_k \varphi_k (g_{kk} - 1/2) \right).$$

In our example $\rho_1 = -T^2 + 2T - 2 + 2T^{-1} - T^{-2}$.

Theorem. ρ_1 is a knot invariant.

Proof: later.

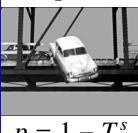
Classical Topologists: Whiskey Tango Foxtrot?



PHOTO NOT AVAILABLE

Cars, Interchanges, and Traffic Counters.

Cars always drive forward. When a car crosses over a bridge it goes through with (algebraic) probability $T^s \sim 1$, but falls off with probability $1 - T^s \sim 0^*$. See also [Jo, LTW].



$$p = 1 - T^s$$

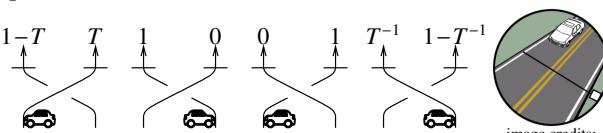


image credits: diamondtraffic.com

* In algebra $x \sim 0$ if for every y in the ideal generated by x , $1 - y$ is invertible.

Preliminaries

This is Rho.nb of <http://drorbn.net/oa22/ap>.

```
Once[<< KnotTheory` ; << Rot.m` ;
```

Loading KnotTheory` version

of February 2, 2020, 10:53:45.2097.

Read more at <http://katlas.org/wiki/KnotTheory>.

Loading Rot.m from <http://drorbn.net/la22/ap>
to compute rotation numbers.

The Program

```
R1[s_, i_, j_] :=  
  s (gji (gj+,j + gj,j+ - gij) - gii (gj,j+ - 1) - 1/2);  
  
Z[K_] := Module[{Cs, ϕ, n, A, s, i, j, k, Δ, G, ρ1},  
  {Cs, ϕ} = Rot[K]; n = Length[Cs];  
  A = IdentityMatrix[2 n + 1];  
  Cases[Cs, {s_, i_, j_}] ↪  
    (A[[{i, j}], {i + 1, j + 1}] += {{-T^s T^s - 1}, {0, -1}})];  
  Δ = T^{(-Total[ϕ] - Total[Cs[[All, 1]]]) / 2} Det[A];  
  G = Inverse[A];  
  ρ1 = Sum[n R1 @@ Cs[[k]] - Sum[n ϕ[[k]] (gkk - 1/2)];  
  Factor@  
  {Δ, Δ^2 ρ1 /. α_ + ↪ α + 1 /. gα_, β_ ↪ G[[α, β]]}];
```

The First Few Knots

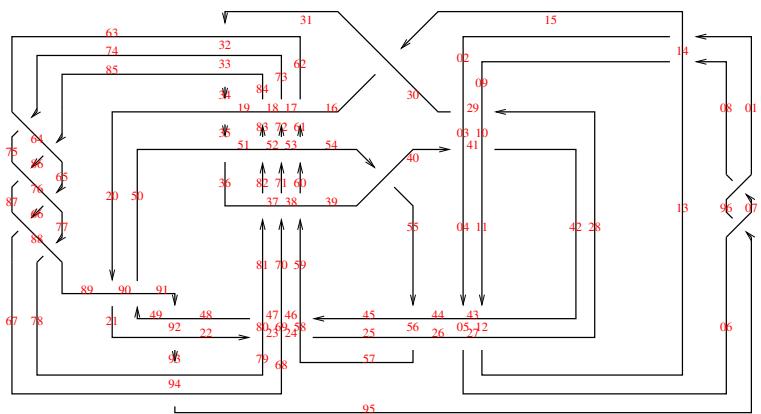
```
TableForm[Table[Join[{K[[1]]K[[2]]}, Z[K]],  
  {K, AllKnots[{3, 6}]}], TableAlignments → Center]
```

3 ₁	$\frac{1-T+T^2}{T}$	$\frac{(-1+T)^2 (1+T^2)}{T^2}$
4 ₁	$-\frac{1-3 T+T^2}{T}$	0
5 ₁	$\frac{1-T+T^2-T^3+T^4}{T^2}$	$\frac{(-1+T)^2 (1+T^2) (2+T^2+2 T^4)}{T^4}$
5 ₂	$\frac{2-3 T+2 T^2}{T}$	$\frac{(-1+T)^2 (5-4 T+5 T^2)}{T^2}$
6 ₁	$-\frac{(2+T) (-1+2 T)}{T}$	$\frac{(-1+T)^2 (1-4 T+T^2)}{T^2}$
6 ₂	$-\frac{1-3 T+3 T^2-3 T^3+T^4}{T^2}$	$\frac{(-1+T)^2 (1-4 T+4 T^2-4 T^3+4 T^4-4 T^5+T^6)}{T^4}$
6 ₃	$\frac{1-3 T+5 T^2-3 T^3+T^4}{T^2}$	0



$$p = 1 - T^s$$

Fast!



Timing@

```
Z[GST48 = EPD[X14,1, X2,29, X3,40, X43,4, X26,5, X6,95,  
X96,7, X13,8, X9,28, X10,41, X42,11, X27,12, X30,15,  
X16,61, X17,72, X18,83, X19,34, X89,20, X21,92,  
X79,22, X68,23, X57,24, X25,56, X62,31, X73,32,  
X84,33, X50,35, X36,81, X37,70, X38,59, X39,54, X44,55,  
X58,45, X69,46, X80,47, X48,91, X90,49, X51,82, X52,71,  
X53,60, X63,74, X64,85, X76,65, X87,66, X67,94,  
X75,86, X88,77, X78,93]]
```

$$\left\{ 170.313, \left\{ -\frac{1}{T^8} (-1 + 2 T - T^2 - T^3 + 2 T^4 - T^5 + T^8) \right. \right. \\ \left. \left. (-1 + T^3 - 2 T^4 + T^5 + T^6 - 2 T^7 + T^8), \frac{1}{T^{16}} \right. \right. \\ \left. \left. (-1 + T)^2 (5 - 18 T + 33 T^2 - 32 T^3 + 2 T^4 + 42 T^5 - 62 T^6 - \right. \right. \\ \left. \left. 8 T^7 + 166 T^8 - 242 T^9 + 108 T^{10} + 132 T^{11} - 226 T^{12} + \right. \right. \\ \left. \left. 148 T^{13} - 11 T^{14} - 36 T^{15} - 11 T^{16} + 148 T^{17} - 226 T^{18} + \right. \right. \\ \left. \left. 132 T^{19} + 108 T^{20} - 242 T^{21} + 166 T^{22} - 8 T^{23} - 62 T^{24} + \right. \right. \\ \left. \left. 42 T^{25} + 2 T^{26} - 32 T^{27} + 33 T^{28} - 18 T^{29} + 5 T^{30} \right) \right\}$$

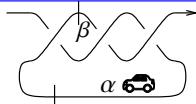
Strong!

```
{NumberOfKnots[{3, 12}],  
Length@  
Union@Table[Z[K], {K, AllKnots[{3, 12}]}],  
Length@  
Union@Table[{HOMFLYPT[K], Kh[K]},  
{K, AllKnots[{3, 12}]}]]}  
{2977, 2882, 2785}
```

So the pair (Δ, ρ_1) attains 2,882 distinct values on the 2,977 prime knots with up to 12 crossings (a deficit of 95), whereas the pair (HOMFLYPT, Khovanov Homology) attains only 2,785 distinct values on the same knots (a deficit of 192).



Theorem. The Green function $g_{\alpha\beta}$ is the reading of a traffic counter at β , if car traffic is injected at α (if $\alpha = \beta$, the counter is *after* the injection point).



Example.

$$\sum_{p \geq 0} (1-T)^p = T^{-1} \quad \begin{array}{c} T^{-1} \\ 1 \end{array} \quad \begin{array}{c} 0 \\ 1 \end{array} \quad \begin{array}{c} 0 \\ 1 \end{array} \quad G = \begin{pmatrix} 1 & T^{-1} & 1 \\ 0 & T^{-1} & 1 \\ 0 & 0 & 1 \end{pmatrix}$$

Proof. Near a crossing c with sign s , incoming upper edge i and incoming lower edge j , both sides satisfy the *g-rules*:

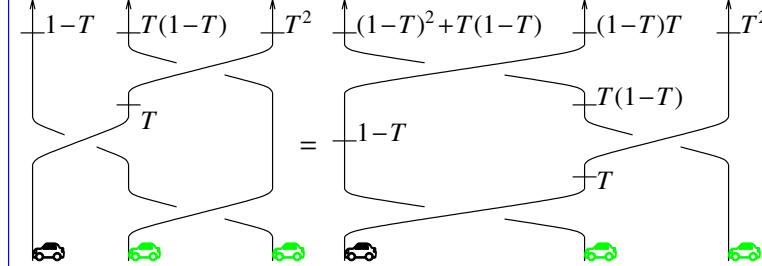
$$g_{i\beta} = \delta_{i\beta} + T^s g_{i+1,\beta} + (1 - T^s) g_{j+1,\beta}, \quad g_{j\beta} = \delta_{j\beta} + g_{j+1,\beta},$$

and always, $g_{\alpha,2n+1} = 1$: use common sense and $AG = I (= GA)$.

Bonus. Near c , both sides satisfy the further *g-rules*:

$$g_{ai} = T^{-s}(g_{\alpha,i+1} - \delta_{\alpha,i+1}), \quad g_{aj} = g_{\alpha,j+1} - (1 - T^s)g_{ai} - \delta_{\alpha,j+1}.$$

Invariance of ρ_1 . We start with the hardest, Reidemeister 3:



⇒ Overall traffic patterns are unaffected by Reid3!

⇒ Green's $g_{\alpha\beta}$ is unchanged by Reid3, provided the cars injection site α and the traffic counters β are away.

⇒ Only the contribution from the R_1 terms within the Reid3 move matters, and using *g-rules* the relevant $g_{\alpha\beta}$'s can be pushed outside of the Reid3 area:

$$\delta_{i,j} := \text{If}[i == j, 1, 0];$$

$$\text{gRules}_{s,i,j} :=$$

$$\left\{ \begin{array}{l} \delta_{i\beta} \Rightarrow \delta_{i\beta} + T^s g_{i^+, \beta} + (1 - T^s) g_{j^+, \beta}, \quad g_{j\beta} \Rightarrow \delta_{j\beta} + g_{j^+, \beta}, \\ g_{\alpha,i} \Rightarrow T^{-s} (g_{\alpha,i^+} - \delta_{\alpha,i^+}), \\ g_{\alpha,j} \Rightarrow g_{\alpha,j^+} - (1 - T^s) g_{\alpha i} - \delta_{\alpha,j^+} \end{array} \right.$$

$$\text{lhs} = R_1[1, j, k] + R_1[1, i, k^+] + R_1[1, i^+, j^+] //.$$

$$\text{gRules}_{1,j,k} \cup \text{gRules}_{1,i,k^+} \cup \text{gRules}_{1,i^+,j^+};$$

$$\text{rhs} = R_1[1, i, j] + R_1[1, i^+, k] + R_1[1, j^+, k^+] //.$$

$$\text{gRules}_{1,i,j} \cup \text{gRules}_{1,i^+,k} \cup \text{gRules}_{1,j^+,k^+};$$

$$\text{Simplify}[\text{lhs} == \text{rhs}]$$

True

Next comes Reid1, where we use results from an earlier example:

$$R_1[1, 2, 1] - 1 (g_{22} - 1/2) /. g_{\alpha,\beta} \Rightarrow \begin{pmatrix} 1 & T^{-1} & 1 \\ 0 & T^{-1} & 1 \\ 0 & 0 & 1 \end{pmatrix} [[\alpha, \beta]]$$

$$\frac{1}{T^2} - \frac{1}{T} - \frac{-1 + \frac{1}{T}}{T} = \textcircled{O}$$

Invariance under the other moves is proven similarly.

Wearing my Topology hat the formula for R_1 , and even the idea to look for R_1 , remain a complete mystery to me.



Wearing my Quantum Algebra hat, I spy a Heisenberg algebra $\mathbb{H} = A\langle p, x \rangle / ([p, x] = 1)$:

$$\text{cars} \leftrightarrow p \quad \text{traffic counters} \leftrightarrow x$$

HEISENBERG

Where did it come from? Consider $g_\epsilon := sl_{2+}^\epsilon := L(y, b, a, x)$ with relations

$$[b, x] = \epsilon x, \quad [b, y] = -\epsilon y, \quad [b, a] = 0,$$

$$[a, x] = x, \quad [a, y] = -y, \quad [x, y] = b + \epsilon a.$$

At invertible ϵ , it is isomorphic to sl_2 plus a central factor, and it can be quantized à la Drinfel'd [Dr] much like sl_2 to get an algebra $QU = A\langle y, b, a, x \rangle$ subject to (with $q = e^{\hbar\epsilon}$):

$$[b, a] = 0, \quad [b, x] = \epsilon x, \quad [b, y] = -\epsilon y, \quad [a, x] = x, \quad [a, y] = -y, \quad xy - qyx = \frac{1 - e^{-\hbar(b+\epsilon a)}}{\hbar}.$$

Now QU has an R -matrix solving Yang-Baxter (meaning Reid3), $R = \sum_{m,n \geq 0} \frac{y^n b^m \otimes (\hbar a)^m (\hbar x)^n}{m! [n]_q!}$, ($[n]_q!$ is a “quantum factorial”)

and so it has an associated “universal quantum invariant” à la Lawrence and Ohtsuki [La, Oh], $Z_\epsilon(K) \in QU$.

Now $QU \cong \mathcal{U}(g_\epsilon)$ (only as algebras!) and $\mathcal{U}(g_\epsilon)$ represents into \mathbb{H} via

$$y \rightarrow -tp - \epsilon \cdot xp^2, \quad b \rightarrow t + \epsilon \cdot xp, \quad a \rightarrow xp, \quad x \rightarrow x,$$

(abstractly, g_ϵ acts on its Verma module

$$\mathcal{U}(g_\epsilon) / (\mathcal{U}(g_\epsilon)(y, a, b - \epsilon a - t)) \cong \mathbb{Q}[x]$$

by differential operators, namely via \mathbb{H}), so R can be pushed to $\mathcal{R} \in \mathbb{H} \otimes \mathbb{H}$.

Everything still makes sense at $\epsilon = 0$ and can be expanded near $\epsilon = 0$ resulting with $\mathcal{R} = \mathcal{R}_0(1 + \epsilon \mathcal{R}_1 + \dots)$, with $\mathcal{R}_0 = e^{t(xp \otimes 1 - x \otimes p)}$ and \mathcal{R}_1 a quartic polynomial in p and x . So p 's and x 's get created along K and need to be pushed around to a standard location (“normal ordering”). This is done using

$$(p \otimes 1)\mathcal{R}_0 = \mathcal{R}_0(T(p \otimes 1) + (1 - T)(1 \otimes p)),$$

$$(1 \otimes p)\mathcal{R}_0 = \mathcal{R}_0(1 \otimes p),$$

and when the dust settles, we get our formulas for ρ_1 . But QU is a quasi-triangular Hopf algebra, and hence ρ_1 is homomorphic. Read more at [BV1, BV2] and hear more at $\omega\epsilon\beta/\text{SolvApp}$, $\omega\epsilon\beta/\text{Dogma}$, $\omega\epsilon\beta/\text{DoPeGDO}$, $\omega\epsilon\beta/\text{FDA}$, $\omega\epsilon\beta/\text{AQDW}$.

Also, we can (and know how to) look at higher powers of ϵ and we can (and more or less know how to) replace sl_2 by arbitrary semi-simple Lie algebra (e.g., [Sch]). So ρ_1 is not alone!

These constructions are very similar to Rozansky-Overbay [Ro1, Ro2, Ro3, Ov] and hence to the “loop expansion” of the Kontsevich integral and the coloured Jones polynomial.

If this all reads like insanity to you, it should (and you haven't seen half of it). Simple things should have simple explanations. Hence, **Homework.** Explain ρ_1 with no reference to quantum voodoo and find it a topology home (large enough to house generalizations!). Make explicit the homomorphic properties of ρ_1 . Use them to do topology!

P.S. As a friend of Δ , ρ_1 gives a genus bound, sometimes better than Δ 's. How much further does this friendship extend?



Schaveling

A Small-Print Page on $\rho_d, d > 1$.

Definition. $\langle f(z_i), h(\zeta_i) \rangle_{\{z_i\}} := f(\partial_{\zeta_i})h|_{\zeta_i=0}$, so $\langle p^2 x^2, \oplus^{g\pi\xi} \rangle = 2g^2$.

Baby Theorem. There exist (non unique) power series $r^\pm(p_1, p_2, x_1, x_2) = \sum_d \epsilon^d r_d^\pm(p_1, p_2, x_1, x_2) \in \mathbb{Q}[T^{\pm 1}, p_1, p_2, x_1, x_2][\epsilon]$ with $\deg r_d^\pm \leq 2d + 2$ (“docile”) such that the power series $Z^b = \sum \rho_d^b \epsilon^d :=$

$$\left\langle \exp\left(\sum_c r^c(p_i, p_j, x_i, x_j)\right), \exp\left(\sum_{\alpha\beta} g_{\alpha\beta} \pi_\alpha \xi_\beta\right) \right\rangle_{\{p_a, x_b\}}$$

is a knot invariant. Beyond the once-and-for-all computation of $g_{\alpha\beta}$ (a matrix inversion), Z^b is computable in $O(n^d)$ operations in the ring $\mathbb{Q}[T^{\pm 1}]$.

(Bnnts are knot diagrams modulo the braid-like Reidemeister moves, but not the cyclic ones).

Theorem. There also exist docile power series $\gamma^\varphi(\bar{p}, \bar{x}) = \sum_d \epsilon^d \gamma_d^\varphi \in \mathbb{Q}[T^{\pm 1}, \bar{p}, \bar{x}][\epsilon]$ such that the power series $Z = \sum \rho_d \epsilon^d :=$

$$\begin{aligned} & \left\langle \exp\left(\sum_c r^c(p_i, p_j, x_i, x_j) + \sum_k \gamma^{\varphi_k}(\bar{p}_k, \bar{x}_k)\right), \right. \\ & \quad \left. \exp\left(\sum_{\alpha\beta} g_{\alpha\beta} (\pi_\alpha + \bar{\pi}_\alpha)(\xi_\beta + \bar{\xi}_\beta) + \sum_\alpha \pi_\alpha \bar{\xi}_\alpha\right) \right\rangle_{\{p_a, \bar{p}_\alpha, x_\beta, \bar{x}_\beta\}} \end{aligned}$$

is a knot invariant, as easily computable as Z^b .

Implementation. Data, then program (with output using the Conway variable $z = \sqrt{T} - 1/\sqrt{T}$), and then a demo. See Rho.nb of $\omega\epsilon\beta/\text{ap}$.

```
V@Y1,_[k_] := ϕ (1/2 - ℙₖ ℙₖ); V@Y2,_[k_] := -ϕ² ℙₖ ℙₖ / 2;
V@Y3,_[k_] := -ϕ³ ℙₖ ℙₖ / 6
V@r1,_[i_, j_] := 
s (-1 + 2 pᵢ xᵢ - 2 pⱼ xⱼ + (-1 + Tˢ) pᵢ pⱼ xᵢ² + (1 - Tˢ) pⱼ xᵢ² - 2 pᵢ pⱼ xᵢ xⱼ + 2 pⱼ² xᵢ xⱼ) / 2
V@r2,_[i_, j_] := 
(-6 pᵢ xᵢ + 6 pⱼ xⱼ - 3 (-1 + 3 T) pᵢ pⱼ xᵢ² + 3 (-1 + 3 T) pⱼ² xᵢ² + 4 (-1 + T) pᵢ² pⱼ xᵢ³ - 
2 (-1 + T) (5 + T) pᵢ pⱼ² xᵢ³ + 2 (-1 + T) (3 + T) pⱼ³ xᵢ³ + 18 pᵢ pⱼ xᵢ xⱼ - 
18 pⱼ² xᵢ xⱼ - 6 pᵢ² pⱼ xᵢ² xⱼ + 6 (2 + T) pᵢ pⱼ² xᵢ² xⱼ - 6 (1 + T) pⱼ³ xᵢ² xⱼ - 
6 pᵢ pⱼ² xᵢ xⱼ + 6 pⱼ³ xᵢ xⱼ) / 12
V@r2,-1[i_, j_] := 
(-6 T² pᵢ xᵢ + 6 T² pⱼ xⱼ + 3 (-3 + T) T pᵢ pⱼ xᵢ² - 3 (-3 + T) T pⱼ² xᵢ² - 
4 (-1 + T) T pᵢ² pⱼ xᵢ³ + 2 (-1 + T) (1 + 5 T) pᵢ pⱼ² xᵢ³ - 2 (-1 + T) (1 + 3 T) pⱼ³ xᵢ³ + 
18 T² pᵢ pⱼ xᵢ xⱼ - 18 T² pⱼ² xᵢ xⱼ - 6 T² pᵢ² pⱼ xᵢ² xⱼ + 6 T (1 + 2 T) pᵢ pⱼ² xᵢ² xⱼ - 
6 T (1 + T) pⱼ³ xᵢ² xⱼ - 6 T² pᵢ pⱼ² xᵢ xⱼ + 6 T² pⱼ³ xᵢ xⱼ) / (12 T²)
```

Z₂[GST48] (* takes a few minutes *)

```
{1 - 4 z² - 61 z⁴ - 207 z⁶ - 296 z⁸ - 210 z¹⁰ - 77 z¹² - 14 z¹⁴ - z¹⁶,
1 + (38 z² + 255 z⁴ + 1696 z⁶ + 16 281 z⁸ + 86 952 z¹⁰ + 259 994 z¹² + 487 372 z¹⁴ + 615 066 z¹⁶ + 543 148 z¹⁸ + 341 714 z²⁰ +
153 722 z²² + 48 983 z²⁴ + 10 776 z²⁶ + 1554 z²⁸ + 132 z³⁰ + 5 z³²) ∈ +
(-8 - 484 z² + 9709 z⁴ + 165 952 z⁶ + 1590 491 z⁸ + 16 256 508 z¹⁰ + 115 341 797 z¹² + 432 685 748 z¹⁴ + 395 838 354 z¹⁶ - 4 017 557 792 z¹⁸ - 23 300 064 167 z²⁰ -
70 082 264 972 z²² - 142 572 271 191 z²⁴ - 209 475 503 700 z²⁶ - 221 616 295 209 z²⁸ - 151 502 648 428 z³⁰ - 23 700 199 243 z³² +
99 462 146 328 z³⁴ + 164 920 463 074 z³⁶ + 162 550 825 432 z³⁸ + 119 164 552 296 z⁴⁰ + 69 153 062 608 z⁴² + 32 547 596 611 z⁴⁴ + 12 541 195 448 z⁴⁶ +
3 961 384 155 z⁴⁸ + 1 021 219 696 z⁵⁰ + 212 773 106 z⁵² + 35 264 208 z⁵⁴ + 4 537 548 z⁵⁶ + 436 600 z⁵⁸ + 29 536 z⁶⁰ + 1252 z⁶² + 25 z⁶⁴) ∈ ²}
```

TableForm[Table[Join[{K[[1]] K[[2]]}, Z₃[K]], {K, AllKnots[{3, 6}]}], TableAlignments → Center] (* takes a few minutes *)

$\begin{array}{ll} 3_1 & 1 + z^2 \\ 4_1 & 1 - z^2 \\ 5_1 & 1 + 3 z^2 + z^4 \\ 5_2 & 1 - 2 z^2 \\ 6_1 & 1 - 2 z^2 \\ 6_2 & 1 - z^2 - z^4 \\ 6_3 & 1 + z^2 + z^4 \end{array}$	$\begin{array}{l} 1 + (2 z^2 + z^4) \in + (2 - 4 z^2 + 3 z^4 + 4 z^6 + z^8) \in^2 + (-12 + 74 z^2 - 27 z^4 - 20 z^6 + 8 z^8 + 6 z^{10} + z^{12}) \in^3 \\ \quad 1 + (-2 + 2 z^2) \in^2 \\ 1 + (6 z^2 + 5 z^4) \in + (4 - 20 z^2 + 43 z^4 + 64 z^6 + 26 z^8) \in^2 + (-36 + 498 z^2 - 883 z^4 + 100 z^6 + 816 z^8 + 556 z^{10} + 146 z^{12}) \in^3 \\ \quad 1 + (-2 z^2 + z^4) \in + (-4 + 4 z^2 + 25 z^4 - 8 z^6 + 2 z^8) \in^2 + (12 + 154 z^2 - 223 z^4 - 608 z^6 + 100 z^8 - 52 z^{10} + 10 z^{12}) \in^3 \\ 1 + (-2 z^2 - 3 z^4 + 2 z^6 + z^8) \in + (-2 - 4 z^2 + 29 z^4 + 28 z^6 + 42 z^8 - 8 z^{10} - 2 z^{12} + 4 z^{14} + z^{16}) \in^2 + (12 + 166 z^2 + 155 z^4 - 194 z^6 - 2453 z^8 - 1622 z^{10} - 1967 z^{12} - 258 z^{14} + 49 z^{16} - 30 z^{18} + z^{20} + 6 z^{22} + z^{24}) \in^3 \\ \quad 1 + (2 + 8 z^2 - 16 z^6 - 24 z^8 - 16 z^{10} - 2 z^{12}) \in^2 \end{array}$
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